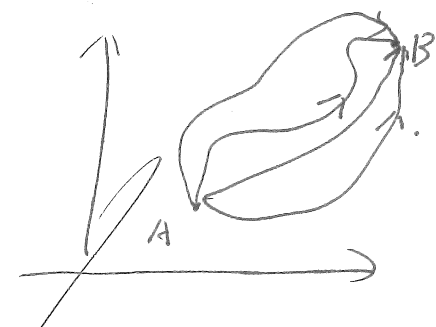


SPUM - pre-lecture. (classical mechanics)

Course idea: To understand the basic concepts of classical mechanics and the fundamental ~~principle~~ mathematics methods of physics. (PDE, ODE, functional, variation, etc.)

1. motion of object.

Consider an object in a configuration space. With all fields covered, it ~~can~~ move from point A \rightarrow B. within a time interval $[t_1, t_2]$.



The trajectory of the motion can be described by some curve line.

$l_1, l_2, \dots, l_n, \dots$ So, How can we know which curve gives us the least action? or say, do to total effective forces do the least effective work?

$$W_{\text{eff}} = \sum_i F_{e,i} \Delta x_i$$

We can use an ~~int~~ functional to estimate ~~this~~ it.

minimize $\rightarrow S = \int_{t_1}^{t_2} \frac{\cancel{W_{\text{eff}}}}{L} dt \Leftrightarrow \int_{t_1}^{t_2} W_{\text{eff}} dt$

where $L = L(q, \dot{q}, t)$. called ~~Lagrangian~~ Lagrangian.

2. principle of least action.

So, how to evaluate this functional?

we could take the variation of it

$$\delta S = \delta \int_{t_1}^{t_2} L dt = \int_{t_1}^{t_2} \delta L dt \quad (\equiv) \quad \int_{t_1}^{t_2} \delta W_{eff} dt$$

Definition: if $\delta S = 0$. then. Lagrangian describes the motion with the least action.

And the L form a PDE

~~$(\delta L = 0)$~~

called Euler-Lagrange Eq.

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} = \frac{\partial L}{\partial q_i} \quad \text{or} \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = 0$$

where $L \equiv T - V$.
 \downarrow \downarrow
 kinetic potential (energy).

0. Mathematical ~~mathematical~~ ~~foundation~~ foundation.

Variation.

$$\delta f(q_1, q_2, \dots, q_n) = \sum_{i=1}^n \frac{\partial f}{\partial q_i} \delta q_i$$

$$\delta \int L dt = \int \delta L dt = \int \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i dt$$

~~multivariable~~
 multivariable. functional.

$$\delta \int L(q_1, q_2, \dots, q_n, \dot{q}_1, \dot{q}_2, \dots, \dot{q}_n, t) dt$$

$$= \int \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i + \sum_{i=1}^n \frac{\partial L}{\partial \dot{q}_i} \delta \dot{q}_i dt.$$

prove:

Consider two functionals. $F(q, \dot{q})$ $\hat{F}(\hat{q}, \hat{\dot{q}})$.

defined as

$$\begin{cases} F(q, \dot{q}) = \int_{t_1}^{t_2} L(q, \dot{q}, t) dt \\ \hat{F}(\hat{q}, \hat{\dot{q}}) = \int_{t_1}^{t_2} L(\hat{q}, \hat{\dot{q}}, t) dt. \end{cases}$$

~~where \hat{q} is another line of trajectory that~~
 for two different paths (q, \dot{q}) $(\hat{q}, \hat{\dot{q}})$.
~~coordinates~~
 paths.

then,

$$F(\hat{q}, \hat{\dot{q}}) - F(q, \dot{q}) = \int_{t_1}^{t_2} L(t, \hat{q}, \hat{\dot{q}}) - L(t, q, \dot{q}) dt.$$

$$= \int_{t_1}^{t_2} dt \left\{ \left[L(t, \hat{q}, \hat{\dot{q}}) - L(t, q, \dot{q}) \right] + \left[\frac{\partial L}{\partial \dot{q}} (\hat{\dot{q}} - \dot{q}) + \frac{\partial L}{\partial \dot{q}^2} (\hat{\dot{q}} - \dot{q})^2 - \frac{\partial L}{\partial q} (q - \hat{q}) - \frac{\partial L}{\partial q^2} (q - \hat{q})^2 \right] + \frac{1}{2} \left[\frac{\partial^2 L}{\partial \dot{q}^2} (\hat{q} - q)^2 + \frac{\partial^2 L}{\partial \dot{q} \partial q} (\hat{q} - q)(\hat{\dot{q}} - \dot{q}) + \frac{\partial^2 L}{\partial \dot{q}^2} (\hat{\dot{q}} - \dot{q})^2 - \frac{\partial^2 L}{\partial q^2} (q - \hat{q})^2 - \frac{\partial^2 L}{\partial q \partial \dot{q}} (q - \hat{q})(\hat{\dot{q}} - \dot{q}) - \frac{\partial^2 L}{\partial \dot{q}^3} (\hat{\dot{q}} - \dot{q})^3 \right] + \dots \right\} dt.$$

denote:

$$\begin{cases} -2 + \hat{z} = 0 \\ (-2 + \hat{z})^2 = 8z^2 \end{cases} \quad \begin{cases} q + \dot{q} = 8\dot{z} \\ (\dot{q} + \dot{z})^2 = 8\dot{z}^2 \end{cases}$$

and $\sqrt{F(\hat{z}, \dot{z})} - F(z, \dot{z}) = \delta F + \delta^2 F + \dots$

So, we get

$$1^{st} \quad \delta F = \int_{t_1}^{t_2} \frac{\partial F}{\partial z} \delta z + \frac{\partial F}{\partial \dot{z}} \delta \dot{z} dt.$$

$$2^{nd} \quad \delta^2 F = \int_{t_1}^{t_2} \frac{\partial^2 F}{\partial z^2} \delta z^2 + \frac{\partial^2 F}{\partial z \partial \dot{z}} \delta z \delta \dot{z} + \frac{\partial^2 F}{\partial \dot{z}^2} \delta \dot{z}^2 dt.$$

⋮

2 E-L eq of motion.

z_i is the general coordinate.

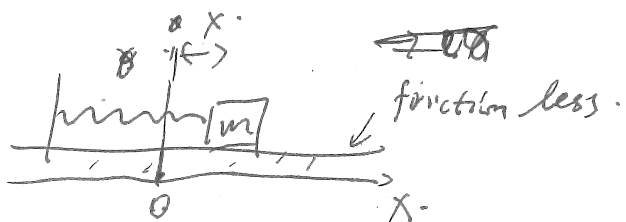
$\dot{z}_i = \frac{dz_i}{dt}$ is the general velocity.

define $\frac{\partial F}{\partial \dot{z}_i} = p_i$ called general momentum.

the equation of motion is

$$\frac{d}{dt} \frac{\partial F}{\partial \dot{z}_i} = \frac{\partial F}{\partial z_i}$$

Ex. Harmonic oscillator.



~~$$L = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$$~~

$$L = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$$

or EOM. (equation of motion). is

$$\frac{d}{dt} \frac{\partial f}{\partial \dot{x}} = \frac{\partial f}{\partial x} \Leftrightarrow m \ddot{x} = -kx$$

$$\underbrace{m}_{\substack{\text{total} \\ \text{mass}}} \underbrace{\ddot{x}}_{\substack{\text{total} \\ \text{accel.}}} = \underbrace{-kx}_{\substack{\text{total} \\ \text{force}}}$$

$$m \ddot{x} = -kx \quad \text{define } \omega = \sqrt{\frac{k}{m}}$$

$$\Rightarrow \ddot{x} + \omega^2 x = 0$$

characteristic eq. $\lambda^2 + \omega^2 = 0$

$$\Rightarrow \lambda = \pm i\omega$$

so. $x(t) = C_1 \sin \omega t + C_2 \cos \omega t$

where C_1, C_2 are some constant.

EX2. ~~pendulum~~ Single pendulum.



$$L = \frac{1}{2} I \dot{\theta}^2 - mgL \cos \theta$$

$V_0 = 0$. E-L eq. should be.

$$\frac{\partial L}{\partial \dot{\theta}} = I \dot{\theta}, \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = I \ddot{\theta} = mL^2 \ddot{\theta}$$

$$\frac{\partial L}{\partial \theta} = mgL \sin \theta$$

$$\Rightarrow mL^2 \ddot{\theta} = mgL \sin \theta$$

$$\Rightarrow \ddot{\theta} = \frac{g}{L} \sin \theta$$

for small damping. $\sin \theta \approx \theta$

$$\ddot{\theta} = \frac{g}{L} \theta \Rightarrow \theta = C_1 e^{\sqrt{\frac{g}{L}} t} + C_2 e^{-\sqrt{\frac{g}{L}} t}$$

- virtual force & work.

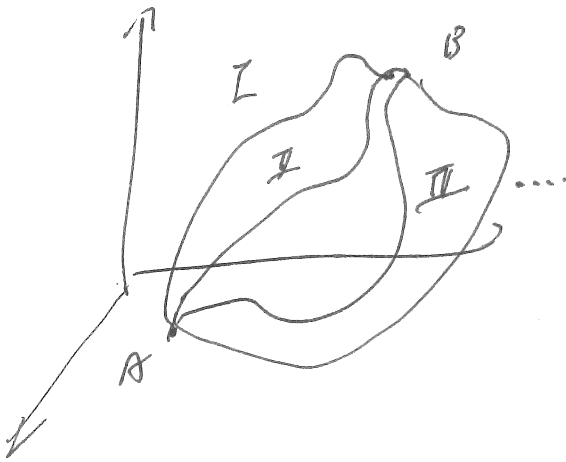
Def: The work of a force acting on a particle as it moves along a displacement is different ~~from~~ for different displacements.
~~The~~ All possible displacements of a particle may follow are called virtual displacement.
~~The~~ The work of a force on a particle along a virtual displacement is known as the virtual work.

I, II, III: virtual displacement.

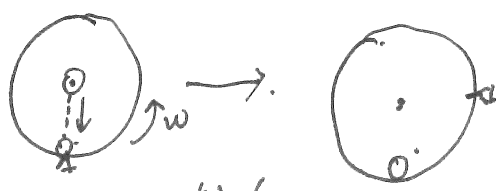
on force F_i doing work δW along it

$$W = \int_{I, II, III} \vec{F}_i \cdot d\vec{l}_i = \sum_i \vec{F}_i \cdot \Delta \vec{q}_i$$

$$\delta W = \sum_i \vec{F}_i \cdot \delta \vec{q}_i \leftarrow \text{virtual force } \delta \vec{q}_i \text{ difference}$$



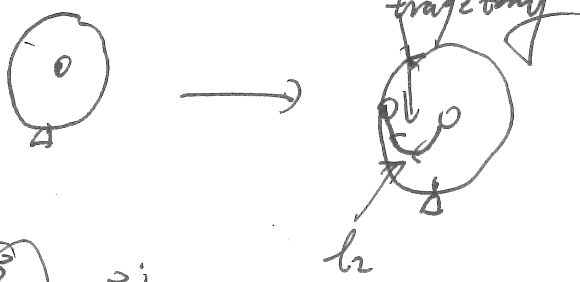
EX1:



In ~~static~~ ^{inertial} stationary reference frame.

In rotating reference frame

the force in ~~static~~ ^{inertial} frame.



$$\vec{F}_I = mg$$

In rotating frame.

$$\vec{F}_R = mg - m \frac{d\vec{\omega}}{dt} \times \vec{R} - 2m\vec{\omega} \times \vec{v} - m\vec{\omega} \times (\vec{\omega} \times \vec{R})$$

$$= \cancel{2}mg - F_{\text{Euler}} - F_{\text{Coriolis}} - F_{\text{centrifugal}}$$

⊗ Nature

$$\int_{l_1} \vec{F}_I \cdot d\vec{l} \equiv \int_{l_2} \vec{F}_R \cdot d\vec{l} \equiv W, \quad \text{and } (\delta W = 0)$$

That means, the total work that done by the fictitious force ~~one~~ is zero. The total virtual work that in the two system are the same. (~~the invariant~~ invariant.)

(~~X~~ If the total ~~work~~ virtual work done by the effective force goes to ~~is zero~~ ~~minimum~~ ~~is zero~~ $(W_e = 0)$ $(\delta W_e = 0)$ then, the ~~that~~ the system is in a dynamic equilibrium.)

~~the, the~~

- D'Alembert's principle.

Def: ~~the~~ the sum of the differences between the forces acting on a system of massive particles and the time derivatives of the momenta of the system itself projected ^{ed} onto any virtual displacement with the constraint of the system is zero.

$$\sum_{i=1}^N (\vec{F}_i - m_i \vec{a}_i) \cdot \delta \vec{r}_i = 0 \quad \Leftrightarrow \quad \sum_{i=1}^N (\vec{F}_i - \dot{\vec{p}}_i) \cdot \delta \vec{r}_i = 0$$

For force and inertia acceleration.

Notice

$$\sum (\mathbf{F}_i - m\mathbf{a}_i) \cdot \delta \mathbf{r}_i = \delta W = 0$$

That ~~the~~ leads to the principle of virtual ~~state~~^{work}
for applied force.

- To the least action principle.

Remember that

$$\delta S = \int_{t_1}^{t_2} \delta W dt \quad \text{if } \delta W = 0.$$

then $\delta S = 0$ that's lead us back to the principle
of least action.